

Outline

- ① Motivation
- ② The Conformal Method
- ③ Bifurcation Theory and unscaled sources
- ④ NON-uniqueness in the extended Conformal Thin Sandwich formulation

MOTIVATION

Pfeiffer + York '05:

(eqn Conformal Thin Sandwich) + (K^0)

= 5 coupled non-linear elliptic PDEs for $(\phi, \tilde{N}, \beta^i)$

Initial Data

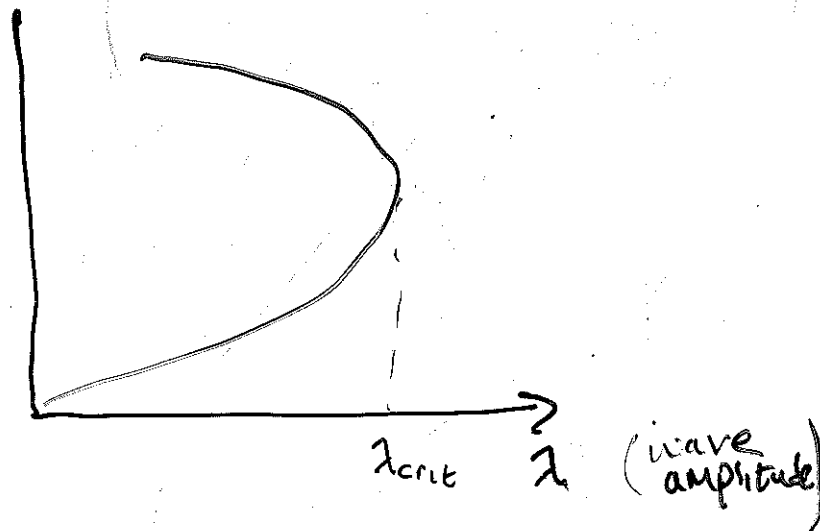
$$\tilde{\gamma}_{ij} = \delta_{ij} + \lambda h_{ij}$$
$$u_{ij} = \lambda \dot{h}_{ij} \quad (\text{tracefree})$$

h_{ij} = Teukolsky wave
 λ = wave amplitude

Results

(max ϕ , min N , max β)

$$\vec{Y} = \vec{Y}_c \pm \sqrt{\epsilon} \vec{W}$$



?

The Conformal Method

$$G_{\mu\nu} = T_{\mu\nu}$$

6 evolution eqns

4 constraint eqns

Constraints

$$R(g) - K_0 K + K^2 = 16\pi \rho$$

$$\nabla_i [k^{ij} - K g^{ij}] = 8\pi j^i$$

choose γ_{ij} s.t. $g_{ij}^{\text{phys}} = \varphi^4 \gamma_{ij}$

$$\Rightarrow R(g) \varphi^5 = R(\gamma) \varphi - 8 \nabla^2 \varphi = \varphi^5 (K_0 K - K^2 + j)$$

$$\left(k_{ij} = A_{ij} + \frac{g_{ij}}{3} K \quad ; \quad \text{take } K = 0 \right. \\ \left. j^i = 0 \right)$$

$$\Rightarrow \boxed{\begin{aligned} \nabla^2 \phi - \frac{R\phi}{8} + \frac{\phi^5}{8} (A_0 A + j_{16\pi}) &= 0 \\ \nabla_i A^{ij} &= 0 \end{aligned}}$$

choose $A_{ij} = \phi^{-10} \hat{A}_{ij}$, $\rho = \psi^{-6} \hat{\rho}$, $K=0$

Linearise LY: (about a solution)

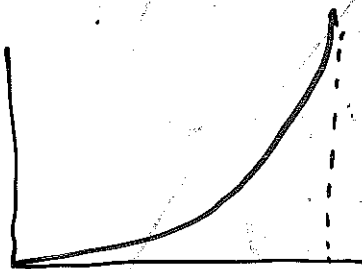
$$(\nabla^2 - f^2) \delta\phi = 0$$

↳ isomorphism (between appropriate spaces)

Imp Fn Thm \Rightarrow there exists unique solutions nearby

Qualitatively $\max \psi$

\sim Pfeiffer + York '05



λ (wave amplitude)

Suppose we don't scale the extrinsic curvature and/or ρ

$$\Rightarrow (\nabla^2 + f) \delta\phi = 0$$

↑
indefinite

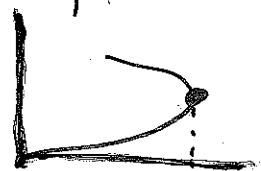
↳ no maximum principle

... linearisation not isomorphic

[note: ^{scale the momentum} many constrained evolutions don't]

Proposal

results of P-Y are due to bifurcation (linearisation of a solution has a kernel) ψ



BASIC ELLIPTIC PDE'S

Linear

$$B u = h(x)$$

, B linear + elliptic

$$x v \quad \int v h = \int v B u = \int (B^* v) u$$

$\therefore B u = h$ is solvable

$\Leftrightarrow \int h v = 0$ for all $v \in \text{Ker } B^*$

non-linear

- Imp. F.n. thm :

$$B u = h(u)$$

, B linear
in our
application

if B is isomorphism
then there exists unique "small"
solutions.

Bifurcation Theory

non-linear eqn $\circledast F : X \rightarrow Y$

$$F = \nabla^2 \phi + \lambda \rho \phi^5 = 0$$

Assume

given : @ $\lambda = \lambda_c$ there's a kernel solution V_0 and critical solution ϕ_c

$$\begin{cases} \nabla^2 \phi_c + \lambda_c \rho \phi_c^5 = 0 \\ \nabla^2 V_0 + 5 \lambda_c \rho \phi_c^4 V_0 = 0 \end{cases}$$

expand in nbd of ϕ_c :

$$\lambda = \lambda_c + \epsilon$$

$$\phi = \phi_c + V$$

$$\circledast \nabla^2 \phi + \lambda \rho \phi^5 \mapsto \nabla^2 (\phi_c + V) + (\lambda_c + \epsilon) \rho (\phi_c + V)^5 = 0$$

$$\Rightarrow \nabla^2 V + \lambda_c \rho \phi_c^5 V = -\epsilon \phi_c^5 \rho - 5\epsilon \phi_c^4 \rho V + 10 \phi_c^3 \rho \lambda_c V^2 + \text{higher order terms}$$

$$\mathcal{B} V$$

$$= \mathcal{R}(V, \epsilon)$$

where \mathcal{B} has 1-D kernel,

$$\mathcal{B} V_0 = 0$$

$$BV = R(V, \epsilon), \quad B: X \rightarrow Y$$

Problem: B has 1-D Kernel v_0

\therefore can't use IFT

Solution: remove the kernel

split Domain $V \in X = \gamma v_0 +$

split Range $R \in Y = dZ +$

$$\begin{pmatrix} u \\ v \\ w \end{pmatrix} \hat{B}$$

(Define $z: Sz \cdot v_0 = 1, \quad Su \cdot z = 0, \quad Sw \cdot v_0 = 0$)

$$BV = R(V, \epsilon)$$

Equivalent to

$$\begin{cases} \textcircled{1} \hat{B}u = R(\gamma v_0 + u, \epsilon) \\ \textcircled{2} Sv_0 R(\gamma v_0 + u, \epsilon) = 0 \end{cases}$$

$\textcircled{1} \hat{B}$ is bijective

\hookrightarrow linearisation of $\textcircled{1}: \hat{B}u = 0$

IFT $\Rightarrow \exists$ solution $u = u(\gamma, \epsilon)$

$\textcircled{2}$ Sub $u(\gamma, \epsilon)$ into $\textcircled{2}$

$$Sv_0 R(\gamma v_0 + u(\gamma, \epsilon), \epsilon) = 0 \quad (\text{i.e. } d=0)$$

$$\Rightarrow \gamma = \gamma(\epsilon)$$

($\textcircled{2}$ is the Lyapunov-Schmidt eqn)

What have we found?

original equation:

$$B V = R(V, \epsilon) \quad \text{s.t.} \quad B V_0 = 0$$

we have a solution

$$V = \underbrace{\gamma(\epsilon) V_0}_{\substack{\text{determined} \\ \text{by} \\ \int V_0 R(\gamma_0 + u, \epsilon) = 0}} + \underbrace{u(\gamma, \epsilon)}_{\text{from IFT}}$$

task find $\gamma = \gamma(\epsilon)$.

recall $R(V, \epsilon) = -\epsilon \phi_c^5 p - 5\epsilon V \phi_c^4 p + 10\phi_c^3 p \gamma V^2 + \dots$

expand V in powers of γ, ϵ

$$V = \underbrace{V_{10}}_{V_0} \gamma + \epsilon V_{01} + \gamma \epsilon V_{11} + \dots$$

sub into $\int R V_0 = 0$:

$$\underbrace{\gamma^2 \int \dots \int V_0}_{\neq 0} + \underbrace{\gamma \epsilon \int \dots \int V_0}_{\neq 0} + \underbrace{\epsilon \int \dots \int V_0}_{\neq 0} \approx 0$$

$$\Rightarrow \gamma = \pm (\text{const}) \sqrt{\epsilon}$$

recall

$$\phi = \phi_c + V = \phi_c + \gamma V_0 + u$$
$$\Rightarrow \boxed{\phi = \phi_c \pm \sqrt{\epsilon} V_0 + \dots} \quad \text{max } \phi$$

Review of Pfeiffer - York '05

Comparison of 2 formulations of Constraints

① CTS : 4 eqn System :

Improved "physical" initial data :

$$(\gamma_{ij}, u_{ij}, K, N)$$

$$(u_{ij} = \text{tracefree } \dot{\gamma}_{ij} / \dot{\gamma}_{ij})$$

Unknowns :

$$(\psi, N, \beta^i)$$

system decouples for $K = \text{const.}$

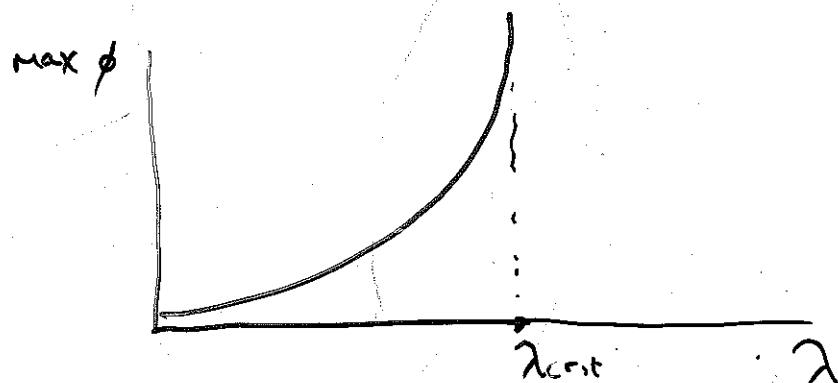
choose

$$\gamma_{ij} = \delta_{ij} + \lambda h_{ij}$$

$$u_{ij} = \dot{\lambda} h_{ij}$$

λ = amplitude of Teukolsky wave

Results



- no solutions for $\lambda > \lambda_{\text{crit}}$
(as expected - Cantor + Brill '81)
and unique solutions for $\lambda < \lambda_{\text{crit}}$
- $E_{\text{ADM}} \rightarrow \infty$ as $\lambda \rightarrow \lambda_{\text{crit}}$

ct2:

2 Extended conformal thin sandwich δ (XCTS)

$$(CTS) + (\dot{K}) = \text{XCTS}$$

$$4 + 1 = 5$$

Initial data $(\sigma_{ij}, u_{ij}, k, \dot{K})$

unknowns (ϕ, β^i, N)

$$\begin{aligned} \nabla^2 \phi - \frac{R\phi}{8} &= -\frac{|\mathbb{L}\beta - u|^2}{32} \frac{\phi^7}{\chi} \\ \nabla^2 \chi - \frac{R\chi}{8} &= \frac{7}{32} |\mathbb{L}\beta - u|^2 \frac{\phi^6}{\chi} \\ \nabla_i \left[\frac{\phi^7}{\chi} (\mathbb{L}\beta^i - u^i) \right] &= 0 \end{aligned}$$

$(\chi := N\phi^7 \text{ chosen } k = \dot{K} = 0)$

• no $k = \text{const.}$ decoupling.

benefit: natural Quasi-Equilibrium

$$u_{ij} = k = \dot{K} = 0$$

But for the Teukolsky wave $(u_{ij} \neq 0)$

$$(\phi, N, \beta) := \vec{Y} = \vec{Y}_c \pm \sqrt{\epsilon} \vec{W}, \quad \epsilon = \delta\lambda$$

$(\max \phi, \min N, \max \beta)$



Can we Apply Lyapunov-Schmidt methods?

- A perturbation problem \circ Assume a solution \vec{y}_c exists whose linearisation has a kernel \vec{V}_0
- need $\text{Ker } \vec{B} = \text{Ker } \vec{B}^*$
 - this is true by examining the linearisation.
- Imp Fn Theorem also valid for systems.

Proceed as before:

Split domain $\vec{V} = \gamma \vec{V}_0 + \vec{u}$
" Range $\vec{R} = \vec{d} \vec{z} + \vec{w}$

($\int \vec{u} \cdot \vec{z} = 0$, $\int \vec{w} \cdot \vec{V}_0 = 0$, ... as before)

Projection:

$$\hat{B} \vec{u} = \vec{R}(\gamma \vec{V}_0 + \vec{u}, \epsilon) \quad (1)$$

$$\int \vec{V}_0 \cdot \vec{R}(\gamma \vec{V}_0 + \vec{u}, \epsilon) = 0 \quad (2)$$

Imp Fn Thm: There exists unique small $\vec{u} = \vec{u}(\gamma, \epsilon)$

Sub into (2) $\Rightarrow \gamma = \gamma(\epsilon)$

\circ $B \vec{V} = R(\vec{V})$ has solution

$$\vec{V} = \gamma(\epsilon) \vec{V}_0 + \vec{u}$$

Taylor expansion of \vec{V} gives $\gamma = \pm \sqrt{\epsilon}$
(assuming very general integrals don't vanish)
Then solution of non-linear problem is

$$\vec{V} = \vec{y}_c \pm \sqrt{\epsilon} \vec{V}_0 + \dots$$

\hookrightarrow Similar to Pflaffer + Vark '05

What have we shown?

- results of P- γ suggest that the linearised system has a kernel at the critical wave amplitude where simple bifurcation effects occur
- Basic non-linear methods can explain complex PDE problems
- should expect similar behaviour in any non-linear elliptic system (if linearisation about exact solution) (has a kernel)

Implications? (with Oliver Rinne)

- a possible reason for the success of free rather than constrained evolutions in some formulations where A_{ij} is not rescaled.
 - $\phi = \dots$ just selects one solution
 - constrained evolution: potential problems on each slice! (if the momentum isn't scaled) ... behaves like S^2
- P- γ results were GLOBAL
 - see also Baumgarte, Ó Murchadha, Pfeffer (106) (analysis of constant density star)
- P- γ used regular initial data
 - much more complicated non-uniqueness using XCTS with isolated horizon boundary conditions (Ansorg, Jaramillo, Luken 106)